

### 3維非定常浮體表面邊界條件

浮體在水面或水中受波運動，勢函數為 $\phi$ ，作自由浮體運動、或被張緊繫留索繫留，靜止時重心位置在 $(\bar{x}_0, \bar{y}_0, \bar{z}_0)$ ，浮心位置在 $(x_b, y_b, z_b)$ ，達平衡狀態時，移至 $(x_0, y_0, z_0)$ ，浮體運動的水平、垂直及回轉角振幅為 $\delta_1$ 、 $\delta_2$ 及 $\delta_3$ （以逆時針方向為正），兩者間關係為

$$\bar{x} = x - (x_0 - \bar{x}_0) + \delta_3(z - \bar{z}_0) - \delta_2(y - \bar{y}_0)$$

$$\bar{y} = y - (y_0 - \bar{y}_0) + \delta_2(x - \bar{x}_0) - \delta_1(z - \bar{z}_0)$$

$$\bar{z} = z - (z_0 - \bar{z}_0) + \delta_1(y - \bar{y}_0) - \delta_3(x - \bar{x}_0)$$

若 $z = \zeta(x, y; t)$ 表示浮體水中部分表面方程式，則由上式得

$$\begin{aligned} \zeta = & -\left[(x_0 - \bar{x}_0) + \delta_2(y - \bar{y}_0) - \delta_3(z - \bar{z}_0)\right] \zeta_x \\ & -\left[(y_0 - \bar{y}_0) + \delta_1(z - \bar{z}_0) - \delta_2(x - \bar{x}_0)\right] \zeta_y \\ & +\left[(z_0 - \bar{z}_0) + \delta_3(x - \bar{x}_0) - \delta_1(y - \bar{y}_0)\right] \zeta_z \end{aligned}$$

上式右邊 $\zeta_x$ ， $\zeta_y$ 及 $\zeta_z$ 為靜止狀態時，浮體水中部份表面S法線在 $x$ ， $y$ ， $z$ 方向分量。由於

$$d\zeta/dt = u\zeta_x + v\zeta_y + w\zeta_z + \zeta_t = 0$$

若流體運動具速度勢 $\Phi(x, y, z; t)$ ，則上式可寫成

$$\Phi_x \zeta_x + \Phi_y \zeta_y - \Phi_z + \zeta_t = 0$$

浮體表面，浮體運動與水粒子運動必須連續得

$$\begin{aligned} \frac{\partial \Phi}{\partial v} = & \frac{\partial x_0}{\partial t} \frac{\partial x}{\partial v} + \frac{\partial y_0}{\partial t} \frac{\partial y}{\partial v} + \frac{\partial z_0}{\partial t} \frac{\partial z}{\partial v} \\ & + \frac{\partial \delta_1}{\partial t} \left[ (z - \bar{z}_0) \frac{\partial y}{\partial v} - (y - \bar{y}_0) \frac{\partial z}{\partial v} \right] \\ & + \frac{\partial \delta_2}{\partial t} \left[ (x - \bar{x}_0) \frac{\partial z}{\partial v} - (z - \bar{z}_0) \frac{\partial x}{\partial v} \right] \\ & + \frac{\partial \delta_3}{\partial t} \left[ (y - \bar{y}_0) \frac{\partial x}{\partial v} - (x - \bar{x}_0) \frac{\partial y}{\partial v} \right] \end{aligned}$$

浮體質量为  $m$ ，通過重心與  $x, y, z$  軸平行慣性力距以  $I_x, I_y, I_z$  表示。作用於浮體壓力除重力、流體壓力  $p$  外，尚有復原力  $R$ ，其在  $x, y, z$  方向分力以  $R_x, R_y, R_z$  表示。復原力所引起對重心力距若以  $N$  表示，其對  $x, y$  及  $z$  軸分力距以  $M_x, M_y, M_z$  表示，則浮體第 1 近似運動方程式為



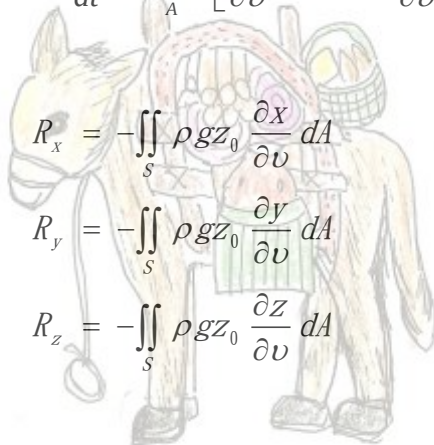
$$\left. \begin{aligned} m \frac{d^2 x_0}{dt^2} &= \iint_A p \frac{\partial x}{\partial v} dA + R_x \\ m \frac{d^2 y_0}{dt^2} &= \iint_A p \frac{\partial y}{\partial v} dA + R_y \\ m \frac{d^2 z_0}{dt^2} &= \iint_A p \frac{\partial z}{\partial v} dA + R_z \end{aligned} \right\}$$



載滿珠寶的駱駝

2011 埃及尼羅河之旅

$$\left. \begin{aligned} I_x \frac{d^2 \delta_1}{dt^2} &= \iint_A p \left[ \frac{\partial z}{\partial v} (y - \bar{y}_0) - \frac{\partial y}{\partial v} (z - \bar{z}_0) \right] dA + M_x \\ I_y \frac{d^2 \delta_2}{dt^2} &= \iint_A p \left[ \frac{\partial x}{\partial v} (z - \bar{z}_0) - \frac{\partial z}{\partial v} (x - \bar{x}_0) \right] dA + M_y \\ I_z \frac{d^2 \delta_3}{dt^2} &= \iint_A p \left[ \frac{\partial y}{\partial v} (x - \bar{x}_0) - \frac{\partial x}{\partial v} (y - \bar{y}_0) \right] dA + M_z \end{aligned} \right\}$$



$$\begin{aligned} R_x &= - \iint_S \rho g z_0 \frac{\partial x}{\partial v} dA \\ R_y &= - \iint_S \rho g z_0 \frac{\partial y}{\partial v} dA \\ R_z &= - \iint_S \rho g z_0 \frac{\partial z}{\partial v} dA \end{aligned}$$



阿拉丁神燈

$$\left. \begin{aligned} M_x &= - \iint_A \rho g \delta_1 (y - \bar{y}_0) \left[ \frac{\partial z}{\partial v} (y - \bar{y}_0) - \frac{\partial y}{\partial v} (z - \bar{z}_0) \right] dA \\ M_y &= - \iint_A \rho g \delta_2 (x - \bar{x}_0) \left[ \frac{\partial x}{\partial v} (z - \bar{z}_0) - \frac{\partial z}{\partial v} (x - \bar{x}_0) \right] dA \\ M_z &= - \iint_A \rho g \delta_3 (z - \bar{z}_0) \left[ \frac{\partial y}{\partial v} (x - \bar{x}_0) - \frac{\partial x}{\partial v} (y - \bar{y}_0) \right] dA \end{aligned} \right\}$$